

Structure + Summary



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- start sph. symm. $*-d.$, glob., nuclear fluid dynamical model, fully self-c. relaxation (r.h.s.) based on Rosenbl, McD, Judd '57
- get corrections for anisotropy decay and equipartition terms
- better agreement with N -body for core collapse, mass seg. (binary heating (not stars))
- generalization to axisym. system poss.
- local \leftrightarrow "straight line" \leftrightarrow orbit average using E, J or E, J_z for axis. rot. syst.
- basis of Monte-Carlo Hénon style
- relation to ultracold plasma (D. Compart) Rydberg atoms

$$k_m = \frac{\lambda}{L}$$

$$\frac{\Delta v}{v} = \frac{2Gm}{pv^2} \lg \frac{\theta}{2}$$

1 (A) Now 6-d Vlasov-Poisson, 1p $f(\vec{x}, \vec{v}, t)$
 sph. symm., collisional term



*-cl: $t_{dyn} \sim \frac{r}{\sigma}$; $t_{rx} \sim \frac{1}{\log \Lambda} t_{dyn}$; $t_{rx} \gg t_{dyn}$
 (in general, not always)

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla f + \underbrace{\vec{v} \cdot \nabla}_{\text{Poisson}} f = \left(\frac{\delta f}{\delta t} \right)_{rx} + \dots + \dots_{diss}$$

Now! (A) L.h.s Vlasov-Poisson

L.h.s. Rosenbluth, 4, 7, 57

Moment equations \leftrightarrow fluid dyn. model
 follow Larson 70, LB + Eggleton 80 Loist+Sp 91

Notation: $\rho = \int_{-\infty}^{\infty} F(\vec{v}) d^3\vec{v} = 2\pi \int_0^{\infty} \int_{-1}^1 d\mu dV V^2 f(V, \mu)$

sph. pol. coord. in v-space
 $\mu = \cos \theta$, $v_r = V\mu$, $v_t^2 = v_\theta^2 + v_\phi^2 = V^2(1-\mu^2)$

Moments: $u\rho = \int \vec{v} F(\vec{v}) d^3\vec{v}$ bulk rad. motion

$\rho_r = \rho \sigma_r^2 = \int (v_r - u)^2 f d^3\vec{v}$ radial p, $\sigma_r \left| \rho = \frac{1}{3}(\rho_r + 2\rho_t) \right.$

$\rho_t = \rho \sigma_t^2 = \int v_\theta^2 f d^3\vec{v} = \int v_\phi^2 f d^3\vec{v}$ tang. p, $\sigma_t \left| e = \frac{3}{2} P \right.$

$F_r = 3\rho_r v_{er} = \int (v_r - u)^3 f d^3\vec{v}$

$F_t = 2\rho_t v_{et} = \int v_r v_\theta^2 f d^3\vec{v} = \int v_r v_\phi^2 f d^3\vec{v}$ thermal energy

radial energy fluxes of rad/tan energy
 net transport velocity of them energy

$F = \frac{1}{2} (F_r + F_t) = \frac{1}{2} (3\rho_r v_{er} + 2\rho_t v_{et}) \stackrel{v_{er} = v_{et}}{=} \frac{5}{2} v_e P$

2

A



- Eqs. of continuity, lin. mom. cons. (Euler) energy (rad + tan)

(Note: r.h.s. = 0!)

Note Euler: $\vec{\sigma} = \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial r} + \frac{\partial \phi}{\partial r} + \frac{1}{\rho} \frac{\partial p_r}{\partial r} + \frac{2 p_a}{\rho r}$

$\frac{p_a}{\rho} = (\sigma_r^2 - \sigma_t^2)$ anisotropy

- Before going to r.h.s.: heat-flux closure

Lynden-Bell + Eggleton 1980:

$F = -\rho \frac{\lambda^2}{T} \frac{\partial \sigma^2}{\partial r}$; λ : MFP, T : $\frac{k_{coll}}{m_{coll}}$ mean free path, $T = \sigma^2$

Use: $\lambda^2 = \frac{1}{k_f^2} = \frac{\sigma^2}{4\pi G \rho}$; $T = T_{rx} = T_{laser}$

Isotr.:

$F = \frac{5}{2} \rho \sigma^2 \cdot \frac{c}{4\pi G \rho T} \frac{\partial \sigma^2}{\partial r} \propto \frac{\rho}{\sigma} \frac{\partial \sigma^2}{\partial r}$

- Generalization to axis. using viscos. $\sim \frac{1}{T}$ (Goodman)

- K_n number related

- external gravity need correction: (Aurore-Seoue, Freitag, Sp 2004)
- $\lambda^2 \rightarrow \lambda^2 \cdot \frac{M(r)}{M_{sh}}$

(only this gets BW disp $\rho \sim r^{-7/4}$)



r.h.s.

Rosenbl., McD., Judd '57, Rosenbluth pot.

$$\left(\frac{\delta h}{\delta t}\right)_{rx} = -4\pi G^2 m_b \ln \Lambda \left[\sum_{i=1}^3 \frac{\partial}{\partial v_i} \left(A(\vec{v}) \frac{\partial h}{\partial v_i} \right) + \frac{1}{2} \sum_{i,j=1}^3 \frac{\partial^2}{\partial v_i \partial v_j} f(\vec{v}) \frac{\partial^2 g}{\partial v_i \partial v_j} \right]$$

$$h(\vec{v}) = (m + m_b) \int \frac{f_b(\vec{v}_b)}{|\vec{v} - \vec{v}_b|} d^3 \vec{v}_b$$

$D(\Delta v_i) A_{chr}$

$$g(\vec{v}) = m_b \int f_b(\vec{v}_b) |\vec{v} - \vec{v}_b| d^3 \vec{v}_b$$

$D(\Delta v_i, \Delta v_j) A_{chr}$

local at \vec{r} , "straight lines"

$$\left[\nabla_v^2 h = -4\pi \left(1 + \frac{m}{m_b}\right) f_b(\vec{v}); \quad \nabla_v^4 g = -8\pi f_b(\vec{v}) \right]$$

NB: RMJ: $d^3 v \leftrightarrow 2\pi V dV d\mu$ (sph. coord. in vel. space)

Other possibility: orbit average, use E, J or E, J_z

l.h.s. = 0 along orbit;

$$r.h.s. = \left(\frac{\delta h}{\delta t}\right)_{rx} = 4\pi G^2 m_b \ln \Lambda \left[\frac{1}{(EJ)^2} \langle \Delta E \rangle \langle \Delta J \rangle + \frac{1}{(EJ)^2} \langle \Delta E^2 \rangle \langle \Delta E^2 \rangle + \frac{1}{(EJ)^2} \langle \Delta E \Delta J \rangle \right]$$

orbit averaged

$$d^3 \vec{v} \leftrightarrow \pi \frac{dE}{v_r} J dJ; \quad v_r = \sqrt{2(E - \phi + \frac{J^2}{m})}$$

$$\mu = \cos \theta; \quad v_r = V \mu; \quad v_+^2 = v_\theta^2 + v_\phi^2 = V^2 (1 - \mu^2)$$

4.

B



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Ansatz following RMB:

$$f(v, \mu) = \sum_{l=0}^{\infty} A_l(v) P_l(\mu)$$

Legendre Polym.

$$A_0(v) = MB$$

$$f_b(v, \mu) = \dots$$

$A_l(v)$ defines l -th moments, eg

$$A_2(v) = \frac{b_r^2 - \sigma_r^2}{3\sigma^2} \frac{v^2}{\sigma^2} A_0(v) \quad \sigma^2 = \frac{1}{3} (b_r^2 - 2A_2)$$

Larson 1970: $f_b = A_0$, f_a up to $l=4$

Louis + Sp 91: $f_b = f$ up to $l=2$

Schneider + Sp 2011: $f_b = f$ up to arbitrary l
for r.h.s. $l=5$ (l.h.s.)

1st non-tr. result: (anisotropy w. non collapse)

$$\left(\frac{dPa}{dT}\right)_{trx} = -\frac{9}{10} \frac{Pa}{T}$$

$$\left(\text{Larson: } -\frac{6}{5} \frac{Pa}{T}\right)$$

not self-cons.

$$Pa/P = \sigma_r^2 - \sigma_t^2$$

$$T = \frac{9}{16\sqrt{\pi}} \frac{G^3}{G^2_{in} \rho \log \Lambda} \quad (\text{"Larson } trx^{\text{u}}")$$



- 2nd non-trivial result: (equipartition)
w. mass segr. (Spt Takahashi's)

two comp. $m, f; m', f'$

$$\left(\frac{\partial p_r}{\partial T}\right)_{rx} = \frac{C_1 + C_2 + C_3 + C_4}{T}$$

$$C_1 = -\frac{3}{4} \rho' \left(\sigma^2 - \frac{m'}{m} \sigma'^2 \right)$$

Spitzer
Eq.

$$C_2 = \frac{1}{5} \rho' \sigma_a^2 \left(1 + \frac{3}{2} \frac{m'}{m} \right) + \frac{1}{5} \rho' \sigma_a'^2 \left(\frac{3}{2} + \frac{m'}{m} \right)$$

$$C_3 = -\frac{3}{10} \rho' \left(\sigma_a^2 + \sigma_a'^2 \right) \frac{\left(1 + \frac{m'}{m} \right)}{\left(1 + \frac{\sigma'^2}{\sigma^2} \right)}$$

$$C_4 = -\frac{27}{140} \rho' \sigma_a^2 \sigma_a'^2 \left(\sigma^2 + \sigma'^2 \right)^{-2} \times \left[\sigma^2 \left(1 + \frac{16}{27} \frac{m'}{m} \right) - \frac{28}{27} \sigma'^2 \left(1 + \frac{39}{28} \frac{m'}{m} \right) \right]$$

Interesting for SIDM?