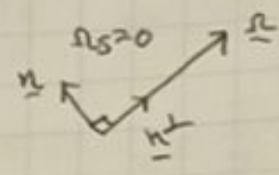


# Nonlinear wave-particle interactions in galaxies (KITP, 28th June 2024)

- Inhomogeneous  $f_0 = f_0(\underline{z})$
- 2D  $\underline{\theta} = (\theta_e, \theta_r)$ ,  $\underline{z} = (z_e, z_r)$
- No collective effects
- Large (nonlinear) external perturbation



Slow-fast angle actions

$$\frac{n_e \Omega_e + n_r \Omega_r - n_e \Omega_p}{\hbar \cdot \underline{n}} = \omega$$

$$\begin{cases} \Omega_s \equiv \underline{n} \cdot \underline{\Omega} & \Omega_p = \underline{n}^\perp \cdot \underline{\Omega} \\ \theta_s = \underline{n} \cdot \underline{\theta} & \theta_p = \underline{n}^\perp \cdot \underline{\theta} \\ J_s = \frac{\underline{n} \cdot \underline{z}}{n^2} & J_p = \frac{\underline{n}^\perp \cdot \underline{z}}{n^2} \end{cases} \quad \left( \begin{array}{l} \underline{\Omega} \equiv (\Omega_e - \Omega_p, \Omega_r) \\ \underline{n} \equiv (n_e, n_r) \\ \underline{n}^\perp \equiv (n_r, -n_e) \end{array} \right)$$

$$S(\underline{\theta}, \underline{z}') = (\underline{n}^\perp \cdot \underline{\theta}) J_p + (\underline{n} \cdot \underline{\theta}) J_s$$

$$J_e = \frac{\partial S}{\partial \theta_e} = n_r J_p + n_e J_s$$

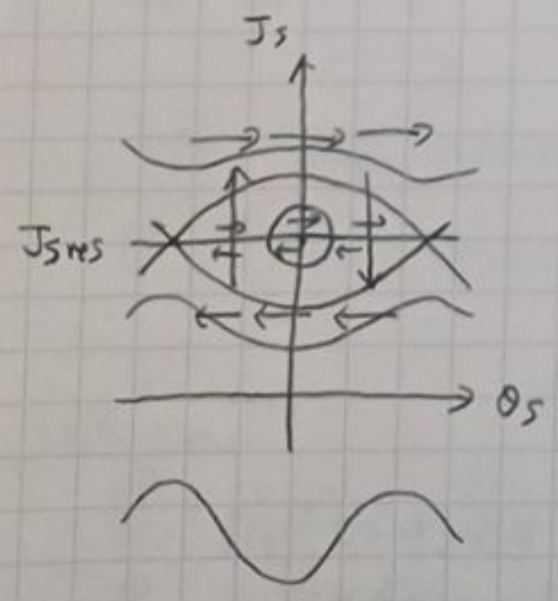
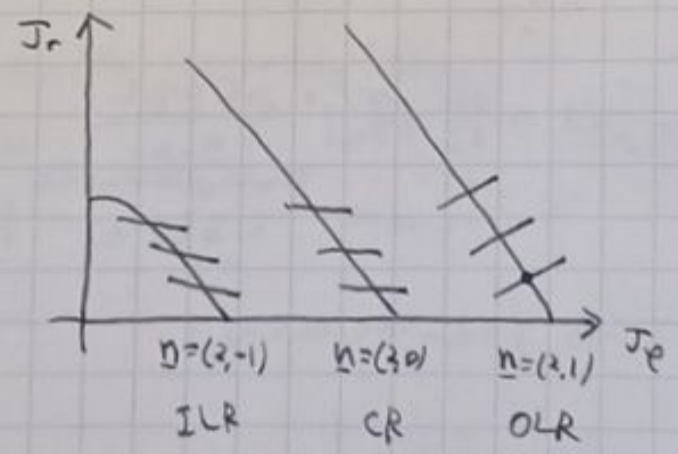
$$J_r = \frac{\partial S}{\partial \theta_r} = -n_e J_p + n_r J_s$$

$$\begin{aligned} n_e J_e + n_r J_r &= (n_e^2 + n_r^2) J_s \rightarrow J_s = \frac{n \cdot \underline{z}}{n^2} \\ n_r J_e - n_e J_r &= (n_r^2 + n_e^2) J_p \rightarrow J_p = \frac{\underline{n}^\perp \cdot \underline{z}}{n^2} \end{aligned}$$

$$H' = H + \frac{\partial S}{\partial t} = H - n_r \Omega_p J_p - n_e \Omega_p J_s = H - \Omega_p (n_r J_p + n_e J_s) = H - \Omega_p J_p$$

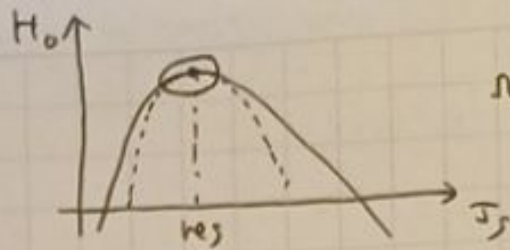
$$\bar{H} = \int \frac{d\theta_p}{2\pi} H$$

$$\dot{z}_p = -\frac{\partial \bar{H}}{\partial \theta_p} = 0$$



$$H = H_0(\mathcal{I}) + H_1(\theta_s, \mathcal{I}, t)$$

$$\approx \frac{1}{2} G (\mathcal{I}_s - \mathcal{I}_{s, \text{res}})^2 - F \cos \theta_s$$



$$\Omega_s = \left. \frac{dH_0}{d\mathcal{I}_s} \right|_{\text{res}} = 0$$

$$G = \left. \frac{\partial^2 H_0}{\partial \mathcal{I}_s^2} \right|_{\text{res}} = \left. \frac{\partial \Omega_s}{\partial \mathcal{I}_s} \right|_{\text{res}} = n_i n_j \frac{\partial \Omega_i}{\partial \mathcal{I}_j} \quad (i, j = \varphi, R)$$

$$\left( \begin{aligned} H_1 &= \sum_m \hat{\Phi}_m(\mathcal{I}, t) e^{i m \cdot \theta} \leftarrow m_s \theta_s + m_f \theta_f \\ &= \sum_{m_s} \hat{\Phi}_{m_s}(\mathcal{I}, t) e^{i m_s \theta_s} \\ &= 2 \sum_{m_s} |\hat{\Phi}_{m_s}| \cos(m_s \theta_s + \varphi_{m_s}) \end{aligned} \right)$$

(Linear theory)

$$f_1(\theta, \mathcal{I}, t) = \sum_n n \cdot \frac{df_0}{d\mathcal{I}} e^{i n \cdot \theta} \int_0^t dt' e^{-i n \cdot \Omega(t-t')} \hat{\Phi}_n(\mathcal{I}, t')$$

time dependence of amplitude

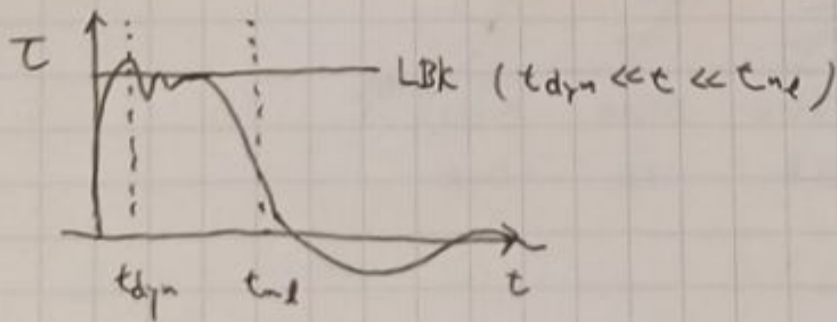
$$= \sum_n n \cdot \frac{df_0}{d\mathcal{I}_s} e^{i n \theta_s} \int_0^t dt' e^{-i n \Omega_s(t-t')} \hat{\Phi}_n(\mathcal{I}, t')$$

$$= \sum_n \frac{df_0}{d\mathcal{I}_s} \hat{\Phi}_n \frac{e^{i n \theta_s} - e^{i(n \theta_s - n \Omega_s t)}}{n \Omega_s}$$

Dynamical friction

$$\text{lin: } \tau = (2\pi)^3 \sum_n n_\varphi \int d\mathcal{I} n \cdot \frac{df_0}{d\mathcal{I}} |\hat{\Phi}_n|^2 \frac{\sin n \cdot \Omega t}{n \cdot \Omega}$$

$t \rightarrow \infty \quad \pi \delta(n \cdot \Omega)$  (Lynden-Bell Kolmogorov formula)



nonlin:  $f(\theta_s, \mathcal{I}_s, t) = f(\theta_0, \mathcal{I}_0, t) = f(\theta_0 - n \Omega t, \mathcal{I}_0, t=0)$

LBK homogeneous:  $M \dot{y}_p = (2\pi)^3 \int d^3 k k |\hat{\Phi}_k|^2 \int d^3 y k \cdot \frac{df_0}{dy} \pi \delta[k \cdot (y - y_p)]$

$$\hat{\Phi}_k = -\frac{\pi k M}{(2\pi)^3 k^2} e^{-i k \cdot y_p t}$$

$$\dot{y}_p = 4\pi G^2 M \ln\left(\frac{k_{\text{max}}}{k_{\text{min}}}\right) \int d^3 y f_0 \frac{y - y_p}{|y - y_p|^3}$$

Chandrasekhar  $\rightarrow$  (HM)

### Moving Resonance

$$\begin{cases} J_s' = J_s - J_{s, \text{res}}(\tau) \\ \theta_s' = \theta_s \end{cases}$$

$$S(\theta_s, J_s', \tau) = \theta_s [J_s' + J_{s, \text{res}}(\tau)]$$

$$H' = H + \frac{dS}{d\tau}$$

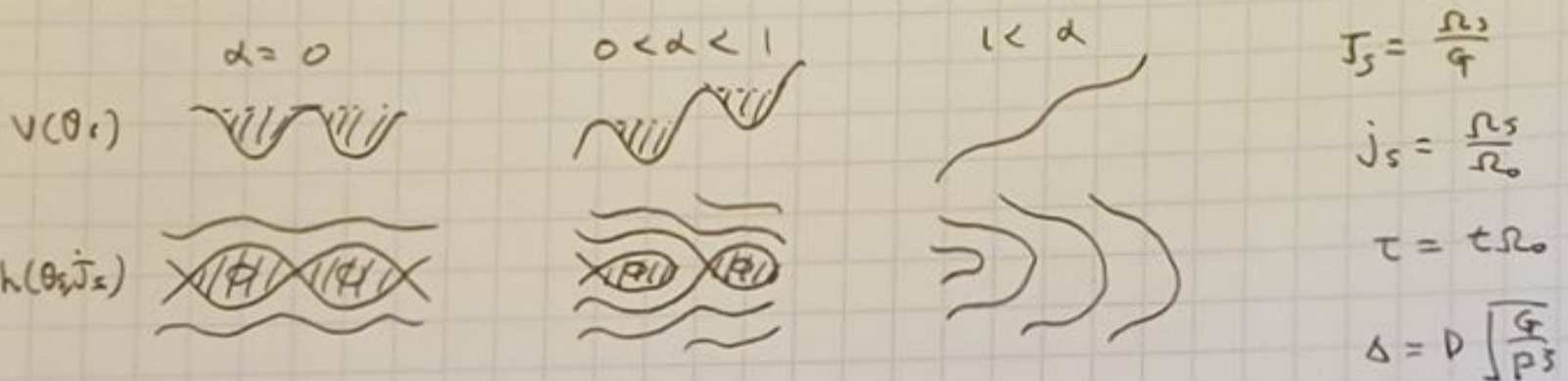
$$= \frac{1}{2} G J_s'^2 - F \cos \theta_s' + J_{s, \text{res}} \theta_s'$$

$$\frac{1}{\tau} \downarrow h = \frac{1}{2} j_s^2 - \underbrace{\omega_s \theta_s + \alpha \theta_s}_{V(\theta_s)}$$

$$J_0 = \sqrt{\frac{F}{G}}, \quad \Omega_0 = \sqrt{F G}$$

$$F = J_0 \Omega_0, \quad G = \frac{\Omega_0}{J_0}$$

$$j_s = \frac{J_s'}{J_0}, \quad \alpha = \frac{J_{s, \text{res}}}{J_0 \Omega_0} \sim \frac{T_{\text{lib}}}{T_{\text{drift}}} = \frac{2}{\pi} \frac{T_{\text{lib}}}{T_{\text{drift}}}$$



### Dynamical feedback

$$\tau_{\text{fb}} = 2\pi |3 \dot{n}_p \sum_n n e \int d^3x \text{sgn}(G) (M_0 - M_e) \delta(n - \Omega)$$

$$M_e = \int_0^{J_{s, \text{res}}} dJ_e f(J_e, J_s, \tau), \quad M_0 = \int_0^{J_{s, \text{res}}} dJ_e f$$

$$\tau_{\text{fb}} = \dot{n}_p (I \dot{n}_p)$$

$$I \dot{n}_p = \tau_{\text{friction}} \tau_{\text{feedback}}$$

$$(I - c) \dot{n}_p = \tau_{\text{friction}}$$

$$M \dot{n}_p = P_{\text{friction}} + P_{\text{feedback}}$$

$$(M - c) \dot{n}_p = P_{\text{friction}}$$

### Diffusion

$$\frac{dF}{d\tau} + [F, H] = D \frac{\partial^2 F}{\partial J_s^2}$$

$$\frac{dF}{d\tau} + j_s \frac{\partial F}{\partial \theta_s} - (\sin \theta_s + \alpha) \frac{\partial F}{\partial J_s} = D \frac{\partial^2 F}{\partial J_s^2}$$

$$- \sin \theta_s \frac{\partial F}{\partial J_s} = \alpha \frac{\partial F}{\partial J_s} + D \frac{\partial^2 F}{\partial J_s^2}$$